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1 Ignition and Flashover of Reduced Scale Compartments with Timber Ceilings

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5 Highlights:

- 6 • Timber
- 7 • Ceiling ignition
- 8 • Flashover

9

10 Abstract:

11 Experiments on reduced-scale compartments with timber ceilings and floors were
12 conducted to investigate the ignition of the ceiling and growth to flashover. Pool fires of
13 various sizes were ignited and measurements were made of temperature in the plume
14 and ceiling jet, and mass of the timber ceiling and floor. The compartment geometry
15 was also systematically varied. The results showed that ignition of the ceiling occurred
16 when the measured temperature at ceiling level was between 328°C and 347°C. This
17 aligned with intermittent flame impingement on the ceiling from the initiating fire.
18 Following ignition of the ceiling, in most cases the fire subsequently spread across the
19 ceiling resulting, ultimately, in ignition of the floor. It was found that duration of
20 preheating of the ceiling (by the ceiling jet) strongly influenced flame spread. In some
21 cases, ignition occurred, but did not result in continuous flame spread across the ceiling
22 – in these cases it was found that the ceiling jet was cooler due to wider geometry of the
23 compartment. The data were compared against existing correlations from the literature,
24 and it was found that existing methods may be used to predict whether ignition of the
25 ceiling may occur and the conditions for flashover.

26 **Keywords:** Timber; Compartment Fires; Ignition; Fire Growth; Heat Transfer

27

28 Nomenclature

| | |
|-------------|---|
| A | Area of pool fire (cm ²) |
| d | Depth of the compartment (mm) |
| h | Height of the compartment (mm) |
| \dot{m} | Burning rate of the fuel prior to the ignition of the ceiling (g/s) |
| \dot{Q} | Heat release rate (HRR) (kW) |
| \dot{Q}_c | Convective portion of the heat release rate(kW) |
| r | Distance along the ceiling from point of flame impingement (mm) |
| $t_{ig,c}$ | Ignition time of the ceiling (s) |
| $t_{ig,f}$ | Ignition time of the floor (s) |

| | |
|-----------------|--|
| Δt | Time difference between ignition of the ceiling and when flames emerged across the entire width of the ceiling (s) |
| Δt_{ig} | Time difference between ignition of the ceiling and the floor (s) |
| T | Temperature of the gas beneath the ceiling, above the centre of the pool fire, prior to ignition of the ceiling ($^{\circ}\text{C}$) |
| T_{ff} | Temperature of the gas beneath the ceiling, in the far field ($^{\circ}\text{C}$) |
| ΔT | Rise in temperature ($^{\circ}\text{C}$) |
| w | Width of the compartment (mm) |
| x | Horizontal distance from point of impingement (mm) |
| z | Distance between the ceiling and the base of the fire (mm) |
| η | Plume temperature correlation coefficient (-) |

29

30 1. Introduction

31 The ignition of combustible ceilings, and subsequent fire growth is of increasing
 32 relevance to fire safety in the built environment due to the increasing use of engineered
 33 timber in buildings [1]. Timber is a combustible material, and it has been shown that
 34 when used as a lining product within a compartment, the contribution of the timber has
 35 the potential to accelerate fire growth [2]. Many jurisdictions therefore place limits on the
 36 use of timber lining materials for some applications. Characterising the hazard posed by
 37 combustible ceilings is therefore essential when designing compartments with exposed
 38 timber ceilings.

39 Timber is a material about which there is a wealth of literature [3,4]. Although there are
 40 relatively few studies which specifically explore ignition and fire spread in the ceiling
 41 configuration, the concepts of ignition theory and the controlling mechanisms are likely
 42 to still apply. The material properties which govern ignition and fire growth are well
 43 established [5,6]. While the critical heat flux is a system parameter, and values obtained
 44 from previous work are unlikely to translate directly to the ceiling configuration, the role
 45 of thermal inertia and a temperature which characterises the onset of pyrolysis are still
 46 likely to be applicable despite the change in orientation.

47 Ignition of the ceiling is typically driven by heat transfer from the buoyant plume from a
 48 fire arising from the moveable fuel load. When the rising buoyant plume impinges on the
 49 ceiling it will transition from rising vertically to flowing radially outwards from the point of
 50 impingement. As the hot gases spread across the ceiling, they will entrain ambient air,
 51 causing the temperature of the gases beneath the ceiling to decrease with distance
 52 from the point of impingement. Therefore, the maximum temperature of the gas beneath
 53 the ceiling and hence the location where the ceiling will pyrolyse first, will occur at the
 54 point of impingement. Buoyant plumes from fires are a well-characterised phenomenon.
 55 The classical description by Morten et al. [7], allows the mass entrainment of an
 56 axisymmetric plume to be described and the temperature of the gas within the plume to
 57 be predicted as a function of height. Behaviour towards the base of the plume (i.e. in

58 and above the flame) was described by McCaffery based on measurements from
59 diffusions flames. This work distinguished between the continuous, intermittent, and
60 plume region [8]. These principles are now firmly established within the regulatory
61 literature and guidance documents that are used by engineers for the design of various
62 engineering systems [9,10].

63 Characterising the ignition of timber ceilings by buoyant plumes within compartments is
64 therefore a problem at the intersection of various topics where there is a wealth of
65 established knowledge. The approach explored here seeks to use existing plume theory
66 to define the thermal environment near the timber ceiling and relate this to the
67 parameters which control the ignition of timber. This aim is to determine appropriate
68 indicators for ignition of the timber ceiling and, once the ceiling has ignited, the
69 subsequent fire development.

70 **1.1. Previous work**

71 The ignition of timber has been studied in various configurations for more than a
72 century. Bartlett [4] summarised the extensive literature on the ignition of timber. In
73 standardised conditions, critical heat fluxes for piloted and unpiloted ignition are
74 approximately 10-13 kW/m² and 25-33 kW/m² respectively [4]. Such data are potentially
75 useful for examining the ignition of timber ceilings. However, it should also be
76 recognised that in the inverted (ceiling) condition, critical heat fluxes and times to
77 ignition may vary compared to samples tested in the horizontal (floor) condition. Shields
78 et al. [11], for example, found that times to ignition increased threefold for samples
79 tested in the ceiling configuration compared to the floor configuration. These results
80 demonstrate that values for critical heat fluxes or times to ignition that were obtained in
81 “floor” configurations cannot necessarily be applied to “ceiling” configurations.
82 Nevertheless, surface temperature of timber remains a useful metric as it is indicative of
83 the rate of pyrolyzate production – and the potential for these products to ignite in the
84 presence of a pilot. The rate of timber pyrolysis is found to rapidly increase between
85 300°C and 500°C which may be a useful criterion for developing this kind of
86 methodology.

87 The temperature rise within a buoyant plume has been described by McCaffrey [8] as
88 follows:

$$\Delta T \propto \left(\frac{z}{\dot{Q}_c^{2/5}} \right)^\eta \quad \text{Equation 1}$$

89 Where z is the vertical distance to the base of the fire, \dot{Q}_c is the convective portion of the
90 heat release rate (HRR), and η is a dimensionless coefficient which varied as a function
91 of height. This correlation was developed for buoyant plumes that were rising free from
92 obstructions. When a rising plume impinges on a ceiling it will be deflected radially; this
93 was characterised by Alpert [12]. Conditions in the ceiling jet were scaled with heat
94 release, height of rise, and radial distance from a central free-field fire. It was found that:

$$\Delta T \propto \left(\frac{\dot{Q}_c/r}{z} \right)^{2/3} \quad \text{Equation 2}$$

95 However, the focus of Alpert's work was on the detection conditions for thermally
96 actuated and smoke actuated detectors rather than a study of ceiling ignition.

97 You and Faeth [13] also studied the heat transfer to a ceiling from an impinging plume.
98 They found that convection dominated the heat transfer on the ceiling, with radiation
99 comprising less than 20% of the total heat transferred. They also observed that the
100 highest rates of heat transfer (60 kW/m²) were in the region of the stagnation point.
101 Kokkala [14,15] conducted a series of experiments to study the heat transfer to and
102 ignition of PMMA ceilings. The experiments featured a ~300 mm disc of fuel onto which
103 a plume from a fixed diameter burner impinged. Heat flux was measured at the surface
104 of the solid using a water-cooled heat flux gauge, and gas phase temperature was
105 measured using thermocouples. This work linked the heat transfer on the ceiling to
106 scaling of the plume correlation (Eq. 1) but found higher heat fluxes than You and
107 Faeth. This was attributed to the nature of the experimental setup in which the source
108 was a gas burner with a fixed diameter (64 mm). Thus, for higher heat release rates
109 (achieved by higher fuel flow rates from the same orifice), the burner tended towards a
110 momentum driven jet flame rather than a buoyancy driven diffusion flame.

111 Kokkala [15] analysed the time to ignition and attempted to correlate this with the
112 inverse square of applied heat flux. However, there was substantial scatter and it was
113 concluded that attempting to fit a line though the data was "unreasonable". In recent
114 years, researchers have expressed a renewed interest in the conditions under which
115 ceiling configurations will ignite. Of particular relevance in relation to this study is work
116 by Nothard et al. [16], and Kotsovinos et al. [17]. Kotsovinos et al. reported the ignition
117 of the ceiling occurred when flames impinged on the soffit. Their data suggests that this
118 occurred when reported gas temperatures were in the region of 300 °C. Nothard et al.
119 linked ignition of the ceiling to the applied heat flux and associated ignition delay time.
120 Again, these data suggest that temperature measurement may provide a useful
121 indicator of potential for ignition – where directly estimating the heat flux or equating this
122 to small scale tests may not be immediately possible.

123 Once combustible ceiling has ignited, there is the potential for the flames to spread.
124 Flame spread on materials in various orientations has been extensively studied (e.g.
125 [18], [19]). In general, the rate of spread is increasingly governed by convective heat
126 transfer as the orientation of a specimen changes from the vertical to ceiling condition.
127 However, although spread on a ceiling is classified as concurrent flame spread, the
128 phenomenon is less well studied in this configuration. Hinkley and Wraight reported a
129 series of experiments on combustible ceiling linings [20]. These experiments involved a
130 burner placed below a combustible ceiling – with the flames from the burner extending
131 along the ceiling. Burning rate of the ceiling was measured, but rate of fire spread was
132 not directly reported. It was noted that the heat transfer from the hot flames and gases
133 would be sufficient to ignite a combustible floor – and thereafter the fire behaviour would
134 likely be dominated by a shortage of air.

135 More recently, authors have studied char patterns from impinging burners, but did not
136 explore flame spread [21]. Notwithstanding the relative paucity of information on flame
137 spread in the ceiling orientation, various authors have noted the importance of the
138 presence of ceiling flames in creating the conditions necessary for flashover. Commonly
139 accepted definitions for the onset of flashover are achieving a heat flux of 20 kW/m² at
140 low level within the compartment; ceiling temperatures of 600 °C; and flaming outside
141 the compartment [5]. In the context of compartment with timber linings, the experiments
142 of Nothard et al. [16] are particularly notable. These were based on a “long”
143 compartment configuration – where the fire was allowed to spread on the combustible
144 floor. The authors observed that ignition of the Cross Laminated Timber (CLT) was
145 followed by a transition from a “travelling” or “spreading” fire, to a “fully involved fire”,
146 where the flame front progressed rapidly through the compartment.

147 From this short review it is evident that:

- 148 • Ignition of timber in the ceiling configuration is influenced by both the orientation
149 (i.e. being in the “ceiling” condition), and also by the local conditions within the
150 impinging plume.
- 151 • The onset of pyrolysis around 300°C may be a useful criterion for the ignition of
152 timber.
- 153 • Temperatures on the centreline of a plume from a diffusion flame roughly follows
154 the trends described by Eq. 1, even when the rise of the plume is interrupted by
155 impingement on a ceiling. Thus, while the HRR of the initiating fire is important,
156 height is the parameter with the strongest influence on the temperature of the
157 gas at the location of impingement.
- 158 • Flame spread on a ceiling is a concurrent flow condition that is likely to be
159 dominated by convective heat transfer from the gases at the ceiling.
- 160 • Once a flame has spread across the ceiling, the heat flux generated by the
161 flames beneath the ceiling is capable of promoting ignition of items at floor level,
162 accelerating the transition to ventilation controlled burning.

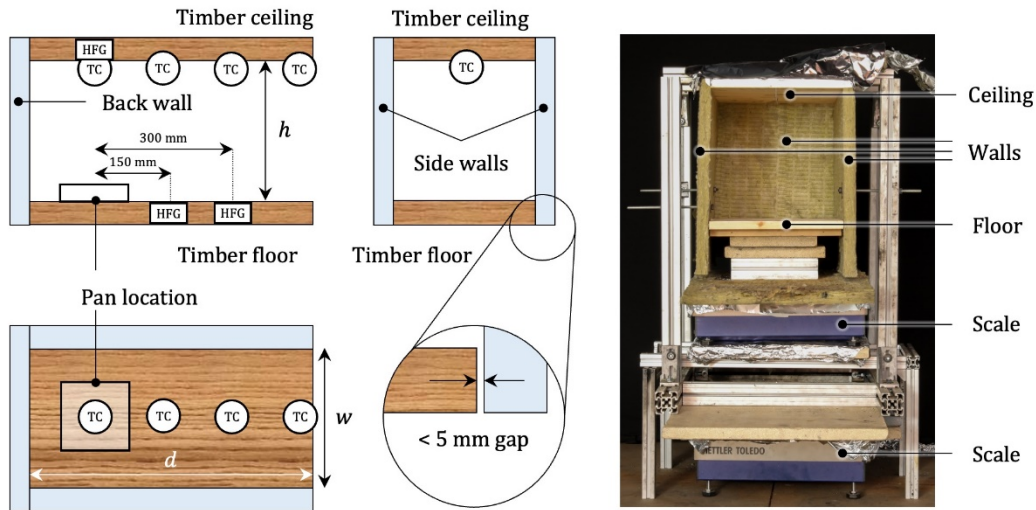
163 These observations allowed the authors to design an experimental study to investigate
164 the applicability of existing work to the ignition and flashover of compartments with
165 timber ceilings.

166 **2. Methodology**

167 As noted previously, the key parameters of interest that feature in previous work are the
168 HRR (\dot{Q}) and the height of rise of the plume (z). The HRR of the initiating fire and
169 separation between the ceiling and base of the fire were systematically varied as part of
170 an experimental study.

171 Pool fires of varying diameters were used to achieve steady fires with different HRRs.
172 The HRR was determined through measurements of the mass and calculation of the
173 mass loss rate (MLR) of the fuel during the experiment. The primary fuel used was
174 paraffin (90%); heptane (10%) was added to the mixture so that the fuel mixture could
175 be easily ignited. Square pans of between 50.4 cm² and 154 cm² were used (giving
176 HRRs of between around 1.5 kW and 5.2 kW). The fuel was ignited using a small flame.

177 An experimental apparatus was designed (Figure 1) which could allow the height of the
 178 compartment to be adjusted. The front of the compartment was open. The pool fire area
 179 (A), depth (d), width (w) and height (h) of the compartment were varied parametrically –
 180 with the primary emphasis being on the variation of \dot{Q} and h .



181
 182 Figure 1 Diagram of the reduced scale compartment presented alongside a
 183 labelled photograph of the experimental setup.

184 A base configuration with a depth of 600 mm, a width of 300 mm and a height of
 185 300 mm was selected. The dimension of 600 mm was selected as this was readily
 186 available from commercial suppliers. From this configuration the compartment height
 187 was increased to 345 mm and 390 mm. The impact of compartment width was also
 188 assessed with a larger configuration; where the depth and width of the compartment
 189 were doubled to 1200 mm and 600 mm respectively. For each of these configurations
 190 the pool area was varied. This approach allowed a scaling analysis to be conducted in
 191 accordance with Eq. 1. A full list of geometries and HRRs for the experiments is
 192 presented as part of Table 1, in the results section.

193 The ceiling and floor of the compartment were 18 mm thick untreated Scandinavian
 194 Redwood Pine furniture board with a density of $470 \pm 10 \text{ kg/m}^3$ and a moisture content
 195 of $8.7 \pm 0.2\%$. The use of a combustible floor was intended to represent other
 196 flammable objects within a compartment. The walls were formed from 25 mm thick
 197 Rockwool Beamclad which was attached to the sides of the ceiling, forming a tight seal.

198 For long or large fires, a 15 mm thick steel plate was placed beneath the fuel pan to act
 199 as a heat sink, so that a steady state pool fire could be maintained (this was preferred
 200 over a water bath which would interfere with the MLR measurements).

201 It should be noted that no downstand was present at the front of the compartment. This
 202 prevented the formation of a smoke layer within the compartment.

203 **2.1 Instrumentation**

204 The instrumentation comprised:

- 205 • Two independent mass balances (of readability 0.1 g) with one supporting the
206 ceiling and walls, and the other supporting the floor and pool fire to allow
207 quantification of the MLR of the pool fire and the time at which ignition of the
208 ceiling and floor occurred.
- 209 • Thermocouples (Inconel sheathed, 1 mm diameter, K type) along the centreline
210 of the compartment in the gas phase (20 mm beneath the ceiling) to allow an
211 estimate of gas temperature above the plume and at 150 mm intervals between
212 the pan and compartment opening. It was assumed that the impact of radiation
213 on the gas temperature measurements were negligible.
- 214 • For selected experiments, water cooled gardon gauges were used to measure
215 the incident heat flux. One located in the ceiling directly above the initiating pool
216 fire and two located in the floor at 150 mm and 300 mm from the centre of the
217 pan.
- 218 • A video camera to allow the time of key events to be identified retrospectively.

219 **2.2 Data Processing**

220 The quasi-steady state HRR of the initiating pool fire preceding ignition of the ceiling
221 was estimated using the change in mass recorded by the mass balance supporting the
222 floor and the fuel pan. This was calculated by averaging the mass loss during the period
223 $t_{ig,c}/2$ to $t_{ig,c}$, where $t_{ig,c}$ is the time to ignition of the ceiling. For experiments where the
224 ceiling did not ignite, the MLR was averaged over half of the time between pan ignition
225 and burnout, centred on the peak gas phase temperature. The HRR was calculated
226 using an assumed combustion efficiency of 0.71 [22], and heat of combustion of 43.1
227 MJ/kg [23]. It was assumed that the change in mass of the bottom scale (Figure 1) prior
228 to ignition of the floor could be attributed to the mass loss of the fuel.

229 The time of the ignition of the ceiling and floor were determined from the MLR recorded
230 by the respective mass balance. The MLR data were smoothed using a 3rd degree
231 Savitzky-Golay filter with a span of 41 data points (equivalent to approx. 40 s). The time
232 of ignition was said to occur when the smoothed MLR began to notably rise above the
233 noise in the MLR. This was found to align with the observed ignition time from
234 experiments and was found to be more definitive in cases where the point of ignition
235 was obscured by the impinging flame.

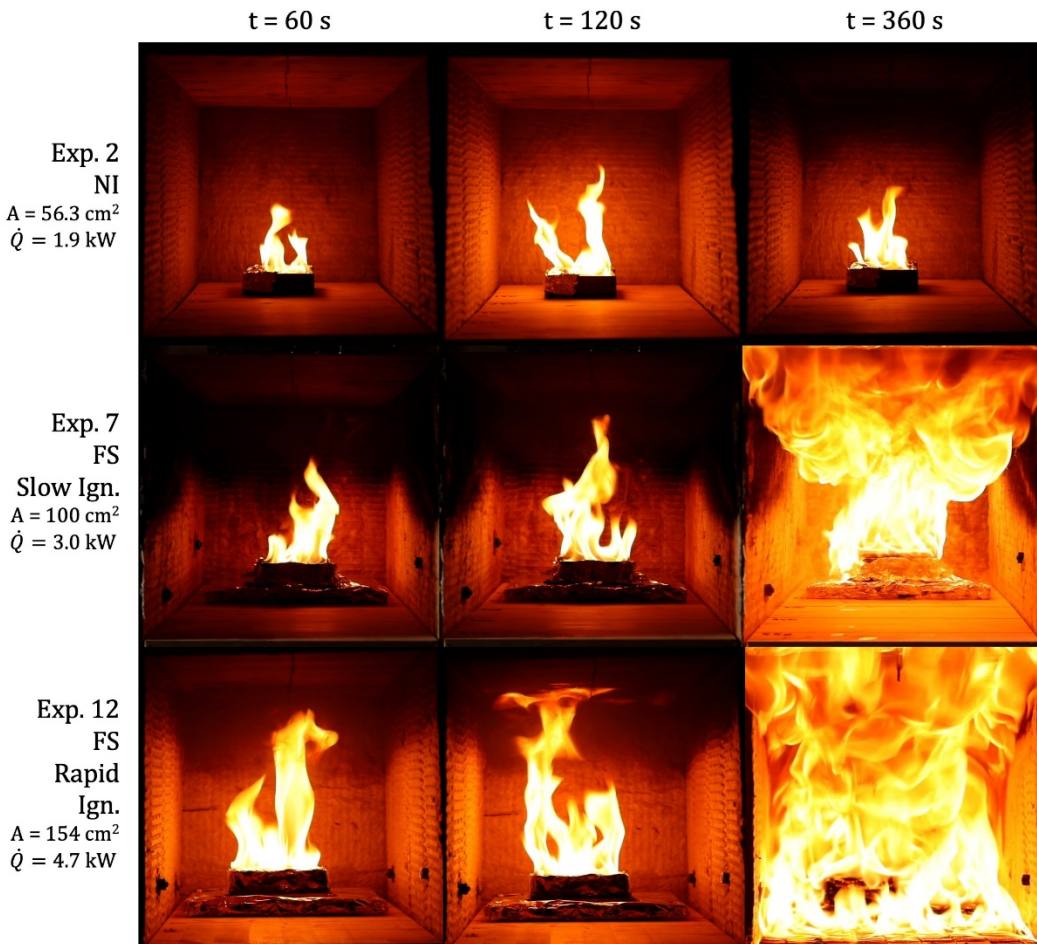
236 The gas phase temperature and heat flux prior to ignition were calculated by averaging
237 their values over the minute preceding ceiling ignition. For experiments where the
238 ceiling did not ignite, the gas phase temperature and heat flux were calculated by
239 averaging their values over a minute centred on the peak gas phase temperature.

240 **3. Results**

241 Twenty-eight experiments were undertaken where a steady state pool fire was
242 achieved. It was observed that the experiments fell into one of three categories. In eight

243 experiments, no ignition of the ceiling occurred (NI). In 16 experiments, ignition of the
 244 ceiling occurred, and this was followed by continuous flame spread across the ceiling
 245 (FS). In four experiments, ignition of the ceiling occurred, but was not followed by
 246 continuous flame spread, these cases are categorised as complex flame spread (CFS).

247 In the 16 FS experiments it was observed that flame spread occurred radially across the
 248 ceiling from the point of ignition to the compartment boundaries. This was followed by
 249 the ignition of the floor around the initiating pool fire, and subsequent flame spread
 250 across the floor from the rear of the compartment to the front, resulting in the
 251 involvement of all combustible surfaces within the compartment. This progression is
 252 shown in Figure 2.



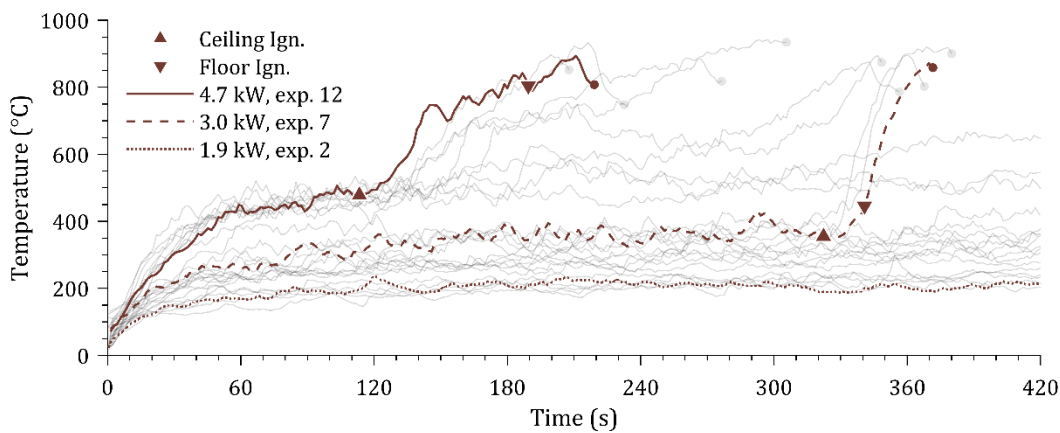
253
 254 Figure 2 Images of three different experiments after 60 s, 120 s, and 360 s;
 255 thermocouple data for the same experiments are presented in Figure 3.

256 Figure 3 shows the temperature measurements at ceiling level directly above the
 257 initiating pool fire for all NI and FS experiments. Three experiments with identical
 258 compartment geometry but different initial HRR have been highlighted to illustrate three
 259 different behaviours: rapid ignition, slow ignition, and no ignition of the ceiling. The time
 260 at which ignition of the ceiling and floor occurred are marked on the highlighted data.

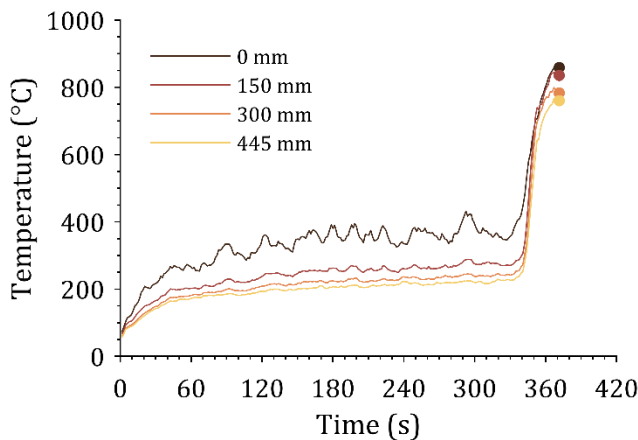
261 The data show that after ignition of the pool fire, the temperature of the gas increased
262 until it achieved a quasi-steady state condition. The ceiling then ignited, leading to a
263 further temperature increase. This was followed by the ignition of the floor. It was
264 observed that the time between the ignition of the ceiling and the floor was greater for
265 larger fires (this is further discussed in the analysis and discussion). This can be seen in
266 the two experiments highlighted in Figure 3. For the experiments where ceiling ignition
267 did not occur, the temperature of the gas remained steady until burnout occurred.

268 The thermocouple data showed that the temperature of the gas beneath the ceiling
269 decreased with distance from the pool fire. This is shown illustratively in Figure 4.

270

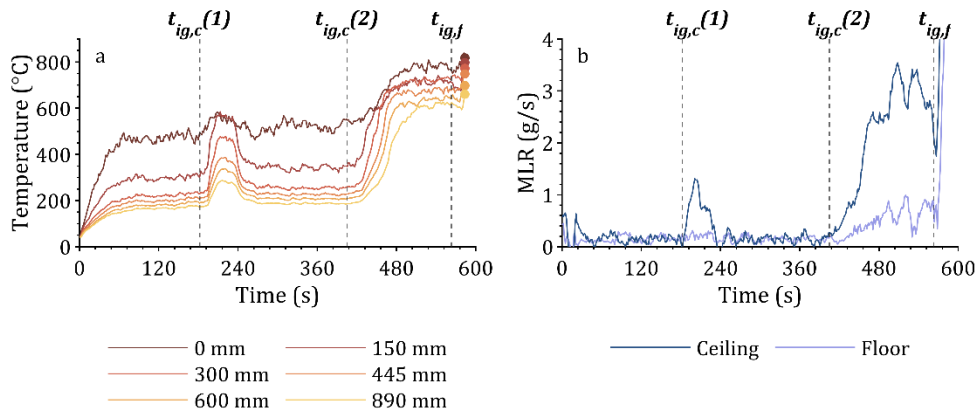


271
272 Figure 3 Temperature beneath the ceiling directly above the initiating fire with three
273 experiments highlighted (Exp. 12, 7 & 2, as presented in Figure 2). For clarity, the data
274 have been curtailed 30 s after the ignition of the floor – represented by a circle.



275
276 Figure 4 Example temperature evolution of the gas beneath the ceiling at different
277 distances from the point of flame impingement(Exp. 7). For clarity, the data have been
278 curtailed 30 s after the ignition of the floor – represented by a circle.

279 In four experiments, the flame spread across the ceiling following ignition was
 280 categorised as “complex”. In these experiments, the flames did not spread continuously
 281 across the entire ceiling following ignition. Instead, flaming was localized around the
 282 point of ignition for a period until the flaming ceased. After further exposure to the
 283 initiating fire, the ceiling of three of these experiments ignited again. It was observed
 284 that, following second ignition of the ceiling, flames spread across the entire ceiling and
 285 the floor ignited, leading to full involvement of all combustible surfaces. An example is
 286 shown in Figure 5 – where the double ignition of the ceiling can be seen in both the
 287 temperature data and the mass loss data. The “complex” behaviours were observed
 288 during the experiments with the widest compartment (600 mm).



289
 290 Figure 5 Example of a) temperature and b) MLR of a CFS exp. (Exp. 25)
 291 The compartment geometry, key parameters and results from each experiment have
 292 been summarised in Table 1. The variable z is the distance between the seat of the pool
 293 fire and the ceiling and accounts for: the thickness of the steel plate (where present),
 294 the thickness of the pan and half of the level of the fuel.

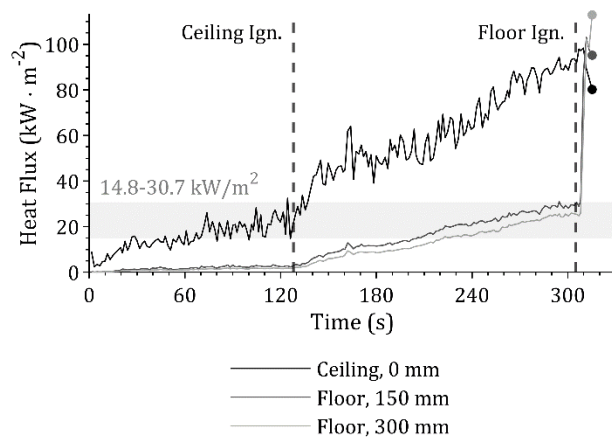
295 Table 1 Experimental matrix and key results

| # | | A cm ² | h mm | d mm | w mm | z mm | \dot{m} g/s | \dot{Q} kW | T °C | $t_{ig,c}$ s | $t_{ig,f}$ s |
|----|----|------------------------|-----------|-----------|-----------|-----------|------------------|-----------------|-----------|-----------------|-----------------|
| 1 | NI | 50.4 | 300 | 600 | 300 | 268 | 0.070 | 2.1 | 300 | - | - |
| 2 | NI | 56.3 | 300 | 600 | 300 | 268 | 0.061 | 1.9 | 245 | - | - |
| 3* | NI | 56.3 | 300 | 600 | 300 | 253 | 0.040 | 1.2 | 224 | - | - |
| 4 | FS | 69.6 | 300 | 600 | 300 | 268 | 0.102 | 3.1 | 360 | 488 | 506 |
| 5 | FS | 100 | 300 | 600 | 300 | 268 | 0.111 | 3.4 | 375 | 325 | 336 |
| 6 | FS | 100 | 300 | 600 | 300 | 268 | 0.094 | 2.9 | 347 | 414 | 423 |
| 7 | FS | 100 | 300 | 600 | 300 | 253 | 0.097 | 3.0 | 372 | 321 | 340 |
| 8* | FS | 100 | 300 | 600 | 300 | 253 | 0.088 | 2.7 | 358 | 322 | 349 |
| 9* | FS | 117 | 300 | 600 | 300 | 243 | 0.128 | 3.9 | 466 | 128 | 305 |
| 10 | FS | 122 | 300 | 600 | 300 | 243 | 0.154 | 4.7 | 469 | 156 | 246 |
| 11 | FS | 122 | 300 | 600 | 300 | 243 | 0.170 | 5.2 | 483 | 141 | 326 |
| 12 | FS | 154 | 300 | 600 | 300 | 253 | 0.153 | 4.7 | 454 | 113 | 189 |

| | | | | | | | | | | | |
|-----|-----|------|-----|------|-----|-----|-------|-----|-----|-----|-----|
| 13 | FS | 154 | 300 | 600 | 300 | 253 | 0.160 | 4.9 | 476 | 122 | 202 |
| 14* | FS | 154 | 300 | 600 | 300 | 253 | 0.100 | 3.1 | 455 | 131 | 177 |
| 15 | NI | 100 | 345 | 600 | 300 | 298 | 0.110 | 3.4 | 328 | - | - |
| 16 | NI | 100 | 345 | 600 | 300 | 298 | 0.097 | 3.0 | 307 | - | - |
| 17 | FS | 117 | 345 | 600 | 300 | 288 | 0.148 | 4.5 | 384 | 535 | 553 |
| 18 | FS | 119 | 345 | 600 | 300 | 288 | 0.104 | 3.2 | 368 | 545 | 552 |
| 19 | NI | 100 | 390 | 600 | 300 | 343 | 0.105 | 3.2 | 280 | - | - |
| 20 | NI | 100 | 390 | 600 | 300 | 343 | 0.102 | 3.1 | 273 | - | - |
| 21 | FS | 117 | 390 | 600 | 300 | 333 | 0.161 | 4.9 | 355 | 681 | 703 |
| 22 | FS | 117 | 390 | 600 | 300 | 333 | 0.151 | 4.6 | 362 | 803 | 806 |
| 23 | NI | 56.3 | 300 | 1200 | 600 | 253 | 0.050 | 1.5 | 228 | - | - |
| 24 | FS | 100 | 300 | 1200 | 600 | 253 | 0.116 | 3.6 | 397 | 465 | 540 |
| 25 | CFS | 117 | 300 | 1200 | 600 | 243 | 0.160 | 4.9 | 477 | 172 | 553 |
| 26 | CFS | 117 | 300 | 1200 | 600 | 243 | 0.165 | 5.0 | 448 | 129 | - |
| 27 | CFS | 117 | 300 | 1200 | 600 | 243 | 0.151 | 4.6 | 503 | 158 | 498 |
| 28 | CFS | 154 | 300 | 1200 | 600 | 253 | 0.130 | 4.0 | 476 | 174 | 317 |

296 *Experiments conducted with heat flux gauges

297 During four experiments (each with the base configuration) the incident heat flux to the
 298 ceiling and the floor was measured (marked in Table 1). An example of the evolution of
 299 the heat flux with time is presented in Figure 6. This plot shows that the heat flux at the
 300 ceiling above the pool fire reaches a quasi-steady state (albeit with substantial
 301 variability) prior to the ignition of the ceiling and the heat flux to the ceiling and floor
 302 increase after ignition of the ceiling. This graph also shows that the ignition of the ceiling
 303 and floor occurred at similar heat fluxes (20.9 kW/m² and 30.7 kW/m² respectively).



304

305 Figure 6 Typical heat fluxes recorded with ignition of the ceiling and floor indicated
 306 (Exp. 9). The range of heat fluxes recorded at ignition of the ceiling and floor are
 307 highlighted.

308 The heat fluxes recorded prior to the ignition of the ceiling and floor are presented in
 309 Table 2. This data show that (unsurprisingly) larger pool fire HRR leads to higher heat
 310 fluxes above the fire. A notable observation from the data is that for Exp. 9 (2.8 kW), the
 311 incident heat flux to the floor was significantly greater than during the other experiments
 312 – these are highlighted in bold in Table 2. The recorded video footage indicates that the
 313 floor in this experiment ignited through autoignition (as opposed to piloted ignition in the
 314 case of the other experiments).

315 Table 2 Recorded heat flux (HF) prior to ignition of the ceiling and floor.

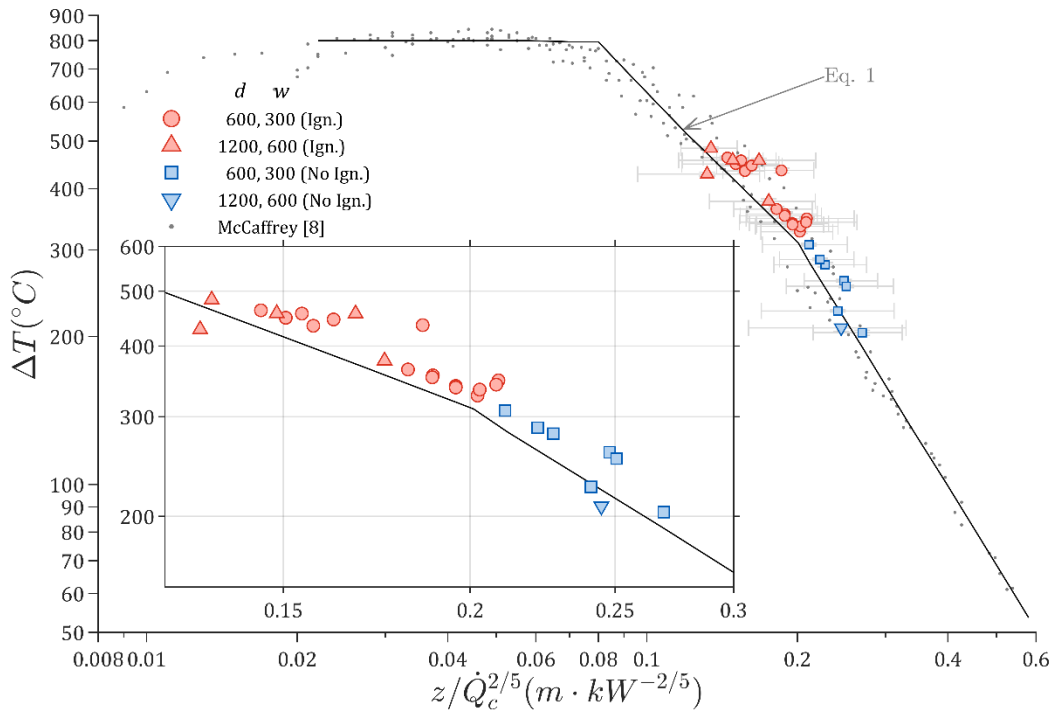
| # | | \dot{Q} kW | $t_{ig,c}$ s | $t_{ig,f}$ s | Δt_{ig} s | Incident Heat flux | | |
|----|----|-----------------|-----------------|-----------------|----------------------|--|--|--------------------|
| | | | | | | Ceiling $t = t_{ig,c}$ $x = 0$ mm kW/m ² | Floor $t = t_{ig,c} / t = t_{ig,f}$ $x = 150$ mm $x = 300$ mm kW/m ² kW/m ² | |
| 3 | NI | 0.91 | - | - | - | 7.83* | 1.04* | 0.66* |
| 8 | FS | 2.1 | 322 | 349 | 27 | 15.2 | 2.2 / 17.8 | 1.44 / 13.2 |
| 9 | FS | 2.8 | 128 | 305 | 177 | 20.9 | 2.6 / 30.7 | 1.69 / 25.4 |
| 14 | FS | 2.5 | 131 | 177 | 46 | 20.7 | 2.7 / 14.8 | 1.96 / 11.6 |

316 *Steady state HF for comparison as no ignition occurred in this case.

317 4. Analysis

318 The initiating fire HRR and compartment geometry were systematically varied, and the
 319 temperature beneath the ceiling where the initiating fire impinged was recorded. This
 320 allowed the results to be compared with the existing fit against Eq. 1 for the variation of
 321 temperature rise with height from the base of the fire and convective portion of the heat
 322 release rate of the fire \dot{Q}_c (assumed to be 70% of the values reported in Table 1 ([5] pg.
 323 136)).

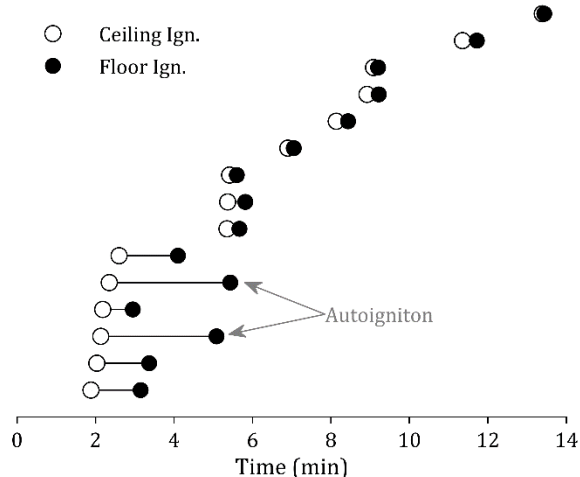
324 These data are presented in Figure 7 alongside the original data. Horizontal error bars
 325 indicate the combined standard error of the height above the pool fire (that will vary as
 326 the fuel is consumed) and standard error of the noise in the HRR. This comparison
 327 shows that the results align with the previous work. The markers are shaded to indicate
 328 whether the ceiling ignited (red) or not (blue). The results show a clear divide between
 329 the experiments where the ceiling ignited and those where it did not. From these results,
 330 the critical ignition temperature lies between 328°C and 347°C. This is within the range
 331 of ignition temperatures of timber reported in literature [24].



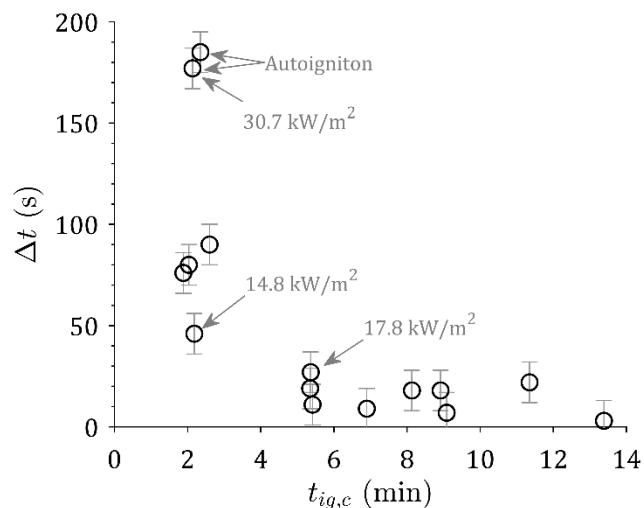
332

333 Figure 7 Comparison between the data in ref. 6, and the data from this study with
 334 gathered data. The depth and width of the compartments are given in millimetres.

335 Following ignition of the ceiling, the flame spread across the ceiling and ignition of the
 336 floor occurred. It was observed that the elapsed time between the ignition of the ceiling
 337 and floor was greater for experiments where the time to ignition of the ceiling was short.
 338 The ignition times of the ceiling and floor are presented in Figure 8 and the variation of
 339 time difference with ignition time is presented in Figure 9. These data were sorted by
 340 ceiling ignition time for presentation in this format. These figures show the importance of
 341 the duration of preheating prior to ignition. The video data showed that for the majority
 342 of the experiments, ignition of the floor was piloted by the fuel burning in the pan.
 343 However, for the two experiments where there was a relatively long elapsed time
 344 between ignition of the ceiling and floor, the floor was observed to ignite through
 345 autoignition (i.e. the pyrolyzate was not observably ignited due to contact with the pool
 346 fire). The heat fluxes recorded in the floor, 150 mm from the centre of the pool fire, at
 347 ignition are marked on Figure 9. This shows that the heat flux was significantly higher
 348 for the two experiments with the larger time difference between ignitions. This
 349 corroborates the notion that the increased elapsed time between ignition of the ceiling
 350 and floor was due to the absence of piloted ignition in these cases.



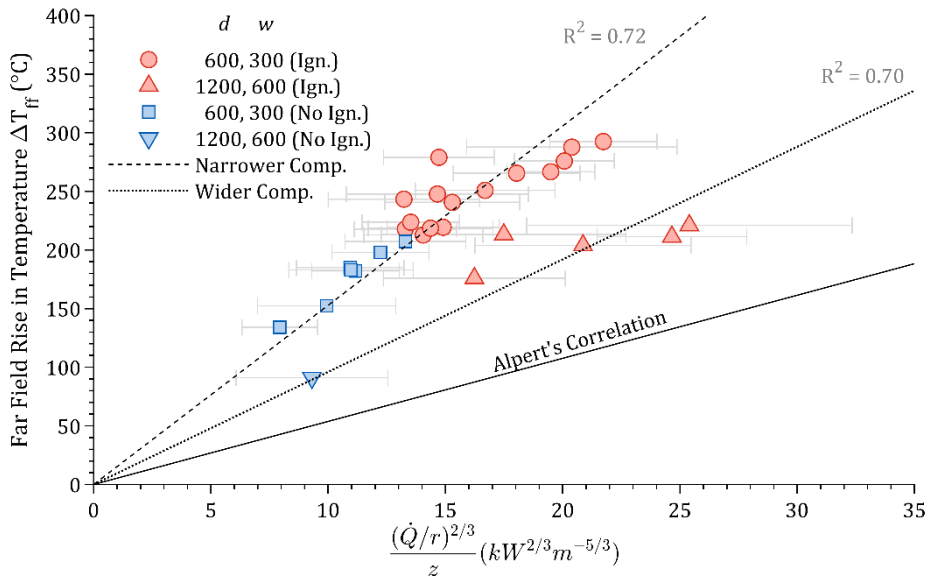
351
 352 **Figure 8** Time to ignition of the ceiling and floor for narrower compartments



353
 354 **Figure 9** Time between ignition of the ceiling and floor against the time to ignition of
 355 the ceiling for narrower compartments. For experiments where the heat flux was
 356 measured, the heat fluxes recorded in the floor, 150 mm from the centre of the pool fire,
 357 at ignition are presented.

358 The preheating of the ceiling is dependent on the condition of the ceiling jet beneath the
 359 ceiling. The temperature of a radially expanding (i.e. unconstrained) ceiling jet is
 360 predicted by Eq. 2. During this experimental programme the ceiling jet was constrained
 361 by the side walls of the compartment – and was not, therefore, free to expand radially in
 362 all directions. Nevertheless, the scaling parameters in Eq. 2 (i.e. initiating fire HRR,
 363 height of rise, and distance from source) remain relevant to the conditions within the
 364 ceiling jet. The temperature, prior to ceiling ignition, 300 mm from the point of
 365 impingement was plotted using these scaling parameters (figure 10). The temperature
 366 measurements are remote from the impinging flames and are referred to as the “far
 367 field” temperatures. This plot shows that Eq. 2 remains useful comparing different

368 temperatures as a function of initial fire HRRs and heights of rise, but that the absolute
 369 value of the temperature (and therefore the constant of proportionality for the scaling
 370 parameter) is significantly influenced by the presence of the compartment boundaries.
 371 The narrower compartment gave the highest temperatures, while the wider
 372 compartment gave lower temperatures. This is due to the reduced entrainment into the
 373 ceiling jet due to the reduced area over which entrainment can occur compared to a
 374 radially expanding ceiling jet. Where the constraint is greater (i.e. the smaller
 375 compartment), then the temperatures are higher. Alpert's correlation (for an
 376 unconstrained expansion of the ceiling jet) is plotted to show the lower bound.



377
 378 Figure 10 Scaling of far field temperature using Alpert's scaling parameters. The
 379 depth and width of the compartments are given in millimetres.

380 **5. Discussion**

381 The temperature increases in the gas phase beneath the ceiling and above a pool fire in
 382 a reduced scale compartment show strong agreement with Eq. 1 and its supporting
 383 data. This echoes previous findings in the context of plume impingement [14] – even
 384 though the original data were developed for an unconfined plume. The results indicate
 385 that ignition of the ceiling can be expected where temperatures in the gas phase are in
 386 the region of 350 °C. Conversely, ignition was not observed where temperatures were
 387 below approximately 330 °C. These values are consistent with ignition temperature
 388 thresholds from the literature. This result demonstrates the potential of existing methods
 389 to remain useful during application to buildings with exposed timber ceilings. The results
 390 suggest that if the gas phase temperature does not exceed 330 °C, the ignition of the
 391 ceiling – and subsequent rapid fire spread through the compartment – may be
 392 prevented.

393 However, it should be noted that the volume of fuel used in the pool fires was intended
 394 to burn for up to 20 minutes (after which time the ceiling was no longer considered
 395 thermally thick). If the duration of the pool fire was extended and thickness of timber

396 increased, then it is possible that some of the experiments where no ignition was
397 observed would have ignited. The potential for substantially delayed ignition of timber
398 has been demonstrated in a set of experiments conducted by Spearpoint and Quintiere,
399 where one sample ignited after 93 minutes of exposure to a constant heat flux of 12
400 kW/m² [24].

401 The presence of compartment boundaries significantly impacts the temperature of the
402 ceiling jet by limiting entrainment, resulting in greater gas temperatures. This was
403 highlighted in the compartments with a greater width where the ceiling jet temperatures
404 were lower (all other factors being equal). In some cases, it was possible for ignition *not*
405 to be followed by immediate fire spread – due to a low level of pre-heating. This
406 illustrates the importance of the entrainment of cool air in the fire spread rate across
407 combustible ceilings.

408 Growth to flashover rapidly follows flame spread on the ceiling. However, this aspect of
409 the experiment was not well controlled as in some cases pyrolysate from the floor was
410 ignited by a pilot (the pool fire), and in some cases pyrolysate from the floor auto-
411 ignited. This indicates that the final event that leads to full involvement of the
412 compartment is stochastic in nature.

413 **6. Conclusion**

414 Evaluation of the hazard posed by timber ceilings in compartments requires knowledge
415 of plume behaviour, conditions for the ignition of timber, flame spread on ceilings, and
416 flashover. Reduced scale experiments have demonstrated that existing knowledge of
417 plume and timber behaviour can be used to determine when ignition of exposed timber
418 ceilings may occur – when subject to heating from a “small” free field pool fire within a
419 compartment. The results show that ceiling ignition, flame spread, and flashover are
420 governed by different processes. It was found that:

- 421 • There was alignment between measured temperatures within the plume above
422 the pool fire and the McCaffrey’s plume correlation [6]. This alignment was found
423 to be a useful method for predicting if ceiling ignition would occur.
- 424 • The preheating of the ceiling and conditions in the smoke layer (or ceiling jet)
425 strongly influenced whether flames spread across the ceiling once ignited.
- 426 • Spread of the fire on the ceiling resulted in a significant increase in heat flux to
427 the floor, which ultimately led to its ignition, and therefore flashover.
- 428 • Once flames had spread across the ceiling, the time delay to ignition of the floor
429 was heavily influenced by stochastic behaviours (i.e. whether the pyrolysis gases
430 were piloted by the pool fire or ignited only due to auto-ignition).

431 The experimental programme allowed the stages of fire development to be
432 systematically examined. These experiments show the utility of existing knowledge in
433 evaluating this hazard, and also the usefulness of relatively small-scale investigations
434 for elucidating the underpinning phenomena that may be observed in larger
435 experiments.

436 However, the scale of these experiments also introduces a limitation. The magnitude of
437 the initiating fire remained constant once established; a fire in a real compartment would
438 be expected to grow, whereas the fire in this experimental programme was constrained
439 to the size of the pan. Similarly, while the scaling analyses are (obviously) intended to
440 be scale independent, these findings have not been corroborated at large scale, and in
441 a design context, should therefore be used with caution.

442

443 7. References

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