



THE UNIVERSITY *of* EDINBURGH

Edinburgh Research Explorer

Polarization independence of laser-induced nucleation in supersaturated aqueous urea solutions

Citation for published version:

Liu, Y, Ward, MR & Alexander, AJ 2017, 'Polarization independence of laser-induced nucleation in supersaturated aqueous urea solutions', *Physical Chemistry Chemical Physics*, vol. 19, no. 5, pp. 3464-3467. <https://doi.org/10.1039/C6CP07997K>

Digital Object Identifier (DOI):

[10.1039/C6CP07997K](https://doi.org/10.1039/C6CP07997K)

Link:

[Link to publication record in Edinburgh Research Explorer](#)

Document Version:

Peer reviewed version

Published In:

Physical Chemistry Chemical Physics

General rights

Copyright for the publications made accessible via the Edinburgh Research Explorer is retained by the author(s) and / or other copyright owners and it is a condition of accessing these publications that users recognise and abide by the legal requirements associated with these rights.

Take down policy

The University of Edinburgh has made every reasonable effort to ensure that Edinburgh Research Explorer content complies with UK legislation. If you believe that the public display of this file breaches copyright please contact openaccess@ed.ac.uk providing details, and we will remove access to the work immediately and investigate your claim.



**Polarization independence of laser-induced nucleation
in supersaturated aqueous urea solutions†**

Yao Liu, Martin R. Ward and Andrew J. Alexander*

*School of Chemistry, University of Edinburgh, David Brewster Road,
Edinburgh, Scotland, EH9 3FJ*

**e-mail: andrew.alexander@ed.ac.uk*

† Electronic supplementary information (ESI) available.

ABSTRACT

In a seminal report on laser-induced nucleation in aqueous supersaturated solutions (*Phys. Rev. Lett.*, 1996, **77**, 3475) it was noted that needle-shaped crystals of urea were aligned with the direction of the electric field of the linearly polarized laser pulse. The results gave rise to a new mechanism for control of crystal nucleation involving alignment of solute molecules (optical Kerr effect) now commonly known as non-photochemical laser-induced nucleation (NPLIN). Recent theoretical and experimental work has cast doubts on the optical Kerr effect mechanism. In the present letter we present results from digital imaging of urea-crystal growth immediately following laser-induced nucleation. Analysis of the data shows no statistically significant correlation between crystal angle and direction of linear polarization. The results overturn a long-held result that has shaped theoretical and experimental studies of NPLIN.

INTRODUCTION

Nucleation is of fundamental scientific interest and has significant economic value. It occurs throughout nature, e.g., in biomineralization, and is used in the production of high-value solid materials, such as pharmaceuticals. A number of methods for inducing nucleation by external perturbations have been employed, for example, electric fields, mechanical shock and ultrasound.¹⁻³ Laser-induced nucleation is a relatively unexplored technique that offers advantages to localizing—temporally and spatially—the point of nucleation.^{4,5} Such control opens up avenues, e.g., to study of the dynamics of nucleation or to control product homogeneity during manufacture of solids.

The discovery of non-photochemical laser-induced nucleation (NPLIN) happened by chance, while Garetz et al. had been looking for second-harmonic generation in aqueous solutions of urea.⁶ They observed nucleation of needle-shaped crystals in supersaturated solutions exposed to pulses of unfocussed, near-infrared laser light (1064 nm, ~20 ns pulse-width, ~2 mm² beam profile, 100–200 pulses at energies of ~100 mJ pulse⁻¹). There are no absorption bands for urea at this wavelength, and multiphoton absorption was discounted due to the relatively low power densities used. Heating of the water (by absorption through weak vibrational overtones) was also considered to be negligible, and if anything, ought to inhibit nucleation by reducing supersaturation. The term *non-photochemical* was introduced to distinguish the effect from nucleation arising from condensation or aggregation due to photochemistry.⁷⁻⁹

Possibly the most significant observation on NPLIN was the correlation between the plane of linear polarization of the light and the direction of the initial urea crystals. To quote Garetz et al.: “*When the laser was vertically polarized, the initial crystallite was observed to be oriented approximately vertically (to within ~5°). When the laser was horizontally polarized, the initial crystallite was oriented approximately horizontally. These correlations were consistently observed for 12.0, 12.5, and 13.0 M solutions*” (the concentration of saturated solution was 10.5 M).⁶ The correlation between crystal axis and polarization of light was strong evidence that NPLIN was a new nucleation phenomenon. The mechanism proposed for NPLIN was based on the optical Kerr effect (OKE), where the most-polarizable axis of the molecule tends to align with the direction of the oscillating electric field of the light. In urea, the direction of the C=O bond is the most-polarizable axis; in the crystal these bonds lie along the [001] direction, which is parallel to the needle axis. In a follow-up study on aqueous supersaturated glycine it was found that the product polymorph could be controlled, with circularly polarized light producing α -glycine and linearly polarized light

producing γ -glycine.¹⁰ It was later noted that complete binary control between the two polymorphs occurred only within a small window of supersaturation (146–155%).¹¹ The influence of the polarization of light on product crystal polymorph through NPLIN has been observed for glycine,¹⁰ L-histidine,¹² carbamazepine,¹³ and sulfathiazole.¹⁴

Several reports have questioned the feasibility of the OKE mechanism for NPLIN. Theoretical considerations indicate that the field strengths involved are too small to align solute molecules sufficiently, even including co-operative effects in solute clusters.^{15, 16} Alternative mechanisms have been put forward.¹⁷ Recent experimental results have pointed towards a mechanism based on heating of trace impurity nanoparticles.^{18, 19} Since the observation of alignment of urea needles along the electric field of the light is central to understanding the mechanism, we have carried out digital imaging of crystal growth during NPLIN of urea.

EXPERIMENTAL METHODS

Urea was obtained from Sigma Aldrich (Sigma Aldrich U5128, 99.0-100.5%) and used without further purification. The solubility of urea (molality) in water at 25 °C is $C_{\text{sat}} = 20.1 \text{ mol kg}^{-1}$.²⁰ Garetz et al. used sample concentrations 11.5–13.0 M (molarity), but it should be noted that molarity depends on the density of the final solution. From tabulated data we estimate that $C = 30.2 \text{ mol kg}^{-1}$ ($S = C / C_{\text{sat}} = 1.5$) corresponds to the 12.5 M solutions.²¹ Supersaturated ($S = 1.5$) aqueous solutions were prepared by dissolving solid urea in ultrapure water (18.2 M Ω cm). Screw-cap glass sample vials (11.7 mm diameter) were cleaned with detergent and rinsed several times, the final rinsing with filtered water [0.22 μm , Millex GP, poly(ether sulfone) membrane]. To ensure equal treatment, all sample vials were prepared and then recrystallized by transferring to a fridge for one day to induce spontaneous nucleation. Supersaturated samples were regenerated by heating in an oven at 50 °C for two days, and the dissolved samples were transferred to another oven with temperature of 25 °C for two days before use.

The laser setup is similar to that described previously.⁴ The light used was obtained from a Q-switched Nd³⁺:YAG laser: either the second harmonic (532 nm, Continuum Surelite II-10, 5.0 ns pulses) or the fundamental (1064 nm, Quantel Brilliant, 5.6 ns pulses). The pulse repetition rate was 10 Hz. The beam was passed through a Glan-laser polarizer to control the transmitted power and to ensure purity of the linear polarization. The beam was passed through a half-wave plate to set the angle of laser polarization with respect to the

vertical direction. A Galilean telescope was used to reduce the beam width, giving final diameters of 4.8 mm (532 nm) or 2.5 mm (1064 nm). The formation of crystals was recorded using a camera (AVT, Stingray F033B) protected by a glass filter (Comar Optics, OG570 or BG39 glass). Each vial was exposed to laser pulses for up to 30 s to induce nucleation; the laser pulses were blocked immediately following nucleation. The vial acts as a cylindrical lens, loosely focusing the beam in the horizontal plane. The refractive index of the solution was estimated to be 1.43.²¹ The peak power density (j_{peak}) was calculated using the area of the beam at the exit of vial. The mean incident laser powers were 0.75 W (532 nm) or 0.31 W (1064 nm), giving peak power densities at the exit of the vials of 200 MW cm⁻² or 270 MW cm⁻², respectively: comparable to those employed by Garetz et al.^{6, 22}

RESULTS AND DISCUSSION

Samples that nucleated showed the quick growth of at least one needle-shaped crystallite (Figure 1), which eventually became a network of needles. The angle (α) between the direction of the electric field of the light and the projection of the initial needle-shaped crystal on the plane of the image (as observed looking along the direction of the beam) was measured from video frames using ImageJ software.²³ It should be noted that α is not the angle between the electric field and the crystal needle axis, but is sufficient to test the correlation noted by Garetz et al. The total number of samples nucleated was 59 at 532 nm and 36 at 1064 nm, respectively. A histogram showing the distribution of the projected angle of the crystal relative to the direction of polarization is shown in Figure 2 (a scatter plot is given in the ESI†). The mean relative angle and standard deviation are given in Table 1. The correlation between the directions appears to be quite random, at both wavelengths.

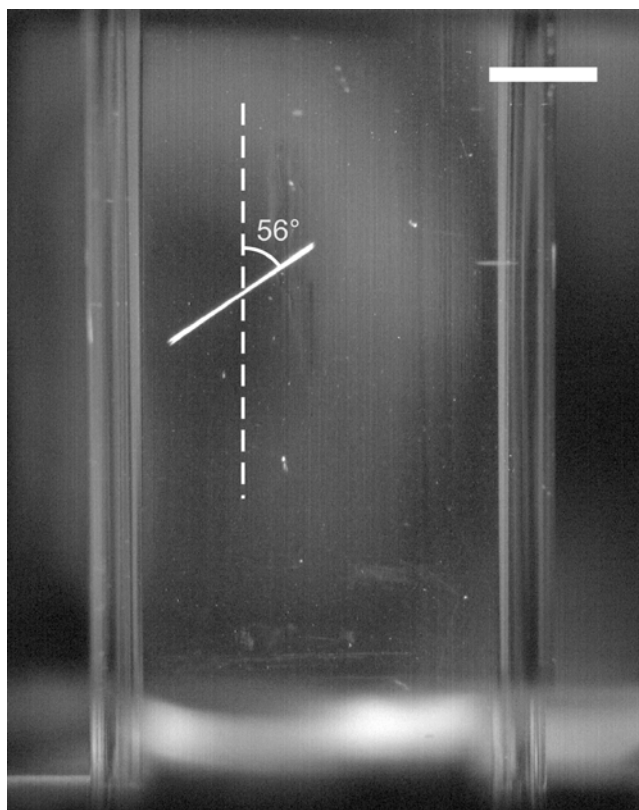


Figure 1. Example image of a urea needle crystal taken moments (0.7 s) after laser-induced nucleation. The laser wavelength was 532 nm and the linear polarization was vertical in the page; the measured projected angle ($\alpha = 56^\circ$) is illustrated. For this example the sample was contained in a square cuvette and the image was taken looking into the oncoming laser beam. The scale bar (top right) represents 3 mm.

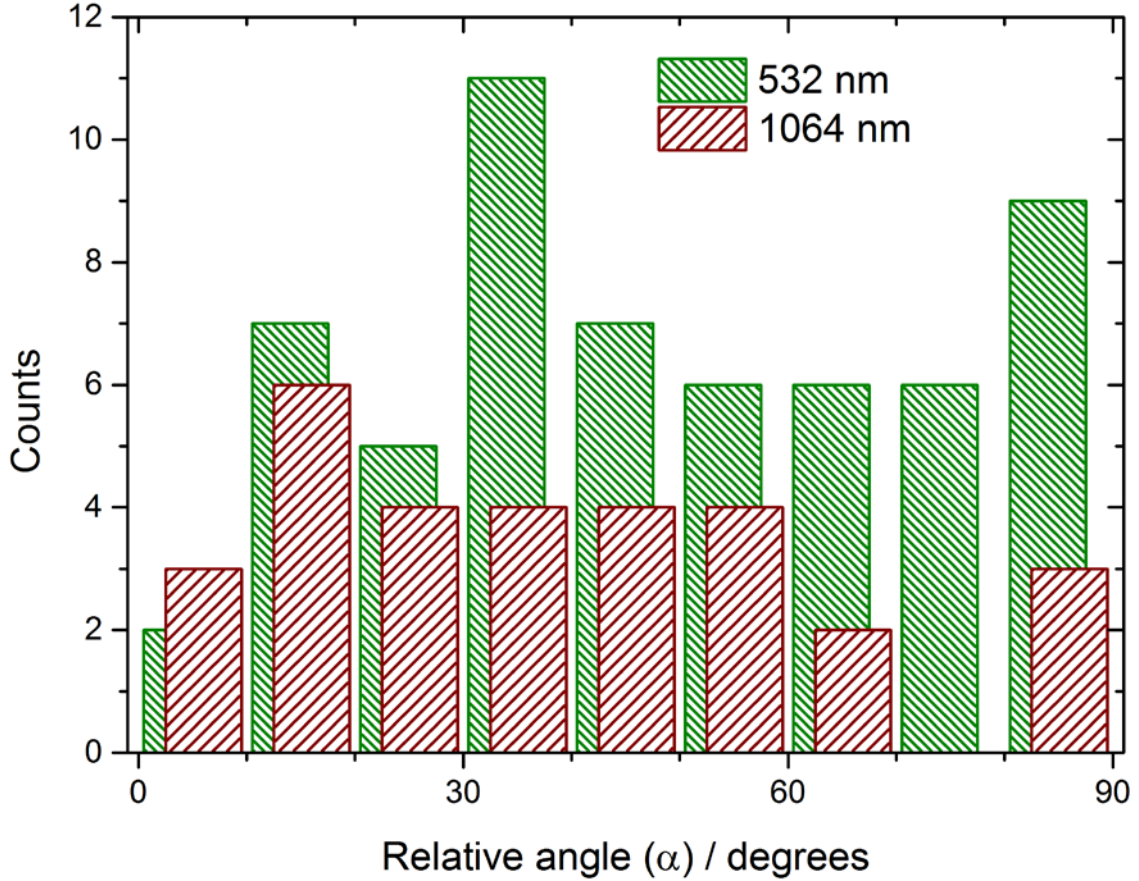


Figure 2. Histogram showing the distribution of the measured projected angle of nascent urea crystals relative to the direction of the electric field of linearly polarized light. The data at 532 nm (green) and 1064 nm (red) show very low correlations. At 1064 nm, 6 samples produced more than 2 crystals: these samples were excluded from the analysis.

To determine if there was a statistically significant correlation in the data, a Kolmogorov-Smirnov (KS) test was performed.²⁴ The data are first arranged in a sequence in increasing numerical order. The KS test compares the cumulative distribution function $F_o(\alpha)$ of the measurements to a continuous model distribution $F_e(\alpha)$. For a sequence of uniformly random angles, we would expect a straight-line model distribution, $F_e(\alpha) = \alpha / (90^\circ)$. The null hypothesis (H_0) was that there was no correlation between laser polarization and crystal direction, i.e., $H_0: F_o(\alpha) = F_e(\alpha)$ for all angles. The alternative hypothesis $H_1: F_o(\alpha) \neq F_e(\alpha)$ for some angles. In order to decide whether the difference between $F_o(\alpha)$ and $F_e(\alpha)$ was significant, we calculated the D statistic, which is the maximum deviation between the distributions:

$$D = \max_{1 \leq i \leq n} [F_o(\alpha_i) - F_e(\alpha_i), F_e(\alpha_i) - F_o(\alpha_{i-1})] \quad (1)$$

where n = number of samples nucleated, and we counted only samples that nucleated a single crystal. The results are summarized in Table 1 (cumulative distributions are given in the ESI†). At a significance level of 0.05, the critical D^* -values corresponding to the data sets at 532 and 1064 nm are $D^*_{59} = 0.174$ and $D^*_{30} = 0.242$, respectively.²⁴ In both cases we find the experimental D -value is less than the critical value, $D < D^*_n$. Therefore, the null hypothesis H_0 should be retained: we find no evidence of a correlation between the observed direction of the needle axis and the direction of the electric field of the linearly polarized light, at a confidence level of 95%. This does not mean that there is no systematic correlation, only that it has not been demonstrated.

wavelength / nm	number of samples (n)	mean relative angle ($\bar{\alpha}$) / deg.	standard deviation of angle (s) / deg.	KS-test statistic (D)	KS-test critical value (D^*)
532	59	48.6	24.2	0.119	0.174
1064	30	37.8	24.1	0.176	0.242

Table 1. Values obtained from analysis of the data. The values for the Kolmogorov–Smirnov (KS) test were used to analyze correlation between the direction of linear polarization and the projected angle of the resulting urea crystal needle. At both wavelengths there appears to be no evidence of a correlation at a confidence level of 95%.

There is no simple explanation for the discrepancy between the original report on urea and the present work. The measurements presented in the original report were quasi-quantitative, with only two polarizations (vertical and horizontal) mentioned. Since the observations were apparently made by eye, it is possible that some delay between nucleation and observation of a sufficiently large needle gave the impression of alignment. Rotation of a growing crystal may be influenced by sedimentation or by convection of the surrounding fluid, which tends to flow toward the crystal.

Apart from NPLIN of urea, Sun et al. reported on alignment of a liquid crystal during cooling through a phase transition.²⁵ It is difficult to draw comparisons to the present experiments, because of significant differences in the experimental setup. Differences include the use of shorter laser pulses (45 ps), exposure to thousands of pulses over an hour of cooling, and the presence of directing polymer surfaces. For some other small-molecule systems, the ellipticity of polarization appears to influence the polymorph

obtained.¹¹⁻¹⁴ In the case of glycine, we note that results from different groups do not agree on the fraction of γ -glycine obtained at different supersaturations.^{11, 26, 27} However, this may be due to different experimental conditions. In the work of Clair et al. the beam passed through both glass–solution and air–solution interface,²⁶ whereas in other studies the beam passed only glass–solution interfaces.^{11, 27}

Studies on NPLIN of carbon dioxide bubbles and on filtration of solutions have suggested that transient heating of an impurity nanoparticle can cause rapid expansion and collapse of a vapor bubble, which triggers nucleation.^{18, 19, 27, 28} The heating mechanism satisfies several observations that the OKE mechanism does not, including the existence of a distinct laser-power threshold, and the fact that some substances do not exhibit NPLIN despite being likely candidates.¹⁹ Absorption and heating would be expected to show little effect of polarization on direction of product nucleation, although further experiments are required to study this systematically.

In summary, the present work has demonstrated that nascent needle-shaped crystals of urea are not aligned with the direction of the electric field as a result of NPLIN. Further experimental and theoretical effort is required to understand the mechanisms underlying laser-induced nucleation.

ACKNOWLEDGEMENTS

Data employed in this study are available via the Edinburgh DataShare repository (doi: 10.7488/ds/1692). We acknowledge funding from the Engineering and Physical Sciences Research Council (EP/I033459/1 and EP/L022397/1) and the EPSRC Center for Innovative Manufacturing in Continuous Manufacturing and Crystallisation (www.cmac.ac.uk).

REFERENCES

1. J. W. Mullin, *Crystallization*, Butterworth-Heinemann, Oxford, 4th edn., 2001.
2. G. Rucroft, D. Hipkiss, T. Ly, N. Maxted and P. W. Cains, *Org. Process Res. Dev.*, 2005, **9**, 923-932.
3. E. Revalor, Z. Hammadi, J.-P. Astier, R. Grossier, E. Garcia, C. Hoff, K. Furuta, T. Okustu, R. Morin and S. Veesler, *J. Cryst. Growth*, 2010, **312**, 939-946.
4. A. J. Alexander and P. J. Camp, *Cryst. Growth Des.*, 2009, **9**, 958-963.
5. C. Duffus, P. J. Camp and A. J. Alexander, *J. Am. Chem. Soc.*, 2009, **131**, 11676-11677.

6. B. A. Garetz, J. E. Aber, N. L. Goddard, R. G. Young and A. S. Myerson, *Phys. Rev. Lett.*, 1996, **77**, 3475-3476.
7. H. Adachi, K. Takano, Y. Hosokawa, T. Inoue, Y. Mori, H. Matsumura, M. Yoshimura, Y. Tsunaka, M. Morikawa, S. Kanaya, H. Masuhara, Y. Kai and T. Sasaki, *Jap. J. Appl. Phys.*, 2003, **42**, L798-L800.
8. N. Nakashima, H. Inoue, M. Sumitani and K. Yoshihara, *J. Chem. Phys.*, 1980, **73**, 4693-4694.
9. T. Okutsu, K. Furuta, M. Terao, H. Hiratsuka, A. Yamano, N. Ferté and S. Veessler, *Cryst. Growth. Des.*, 2005, **5**, 1393-1398.
10. B. A. Garetz, J. Matic and A. S. Myerson, *Phys. Rev. Lett.*, 2002, **89**, 177501.
11. X. Y. Sun, B. A. Garetz and A. S. Myerson, *Cryst. Growth. Des.*, 2006, **6**, 684-689.
12. X. Sun, B. A. Garetz and A. S. Myerson, *Cryst. Growth Des.*, 2008, **8**, 1720-1722.
13. A. Ikni, B. Clair, P. Scoufflaire, S. Veessler, J. M. Gillet, N. El Hassan, F. Dumas and A. Spasojević-de Biré, *Cryst. Growth. Des.*, 2014, **14**, 3286-3299.
14. W. Li, A. Ikni, P. Scoufflaire, X. Shi, N. El Hassan, P. Gémeiner, J.-M. Gillet and A. Spasojević-de Biré, *Cryst. Growth. Des.*, 2016, **16**, 2514-2526.
15. V. Agarwal and B. Peters, *Adv. Chem. Phys.*, 2014, **155**, 97-160.
16. B. C. Knott, M. F. Doherty and B. Peters, *J. Chem. Phys.*, 2011, **134**, 154501.
17. M. Nardone and V. G. Karpov, *Phys. Chem. Chem. Phys.*, 2012, **14**, 13601-13611.
18. M. R. Ward, W. J. Jamieson, C. A. Leckey and A. J. Alexander, *J. Chem. Phys.*, 2015, **142**, 144501.
19. M. R. Ward, A. M. Mackenzie and A. J. Alexander, *Cryst. Growth. Des.*, 2016, doi: 10.1021/acs.cgd.6b00882.
20. A. Seidell, *Solubilities of Organic Compounds*, D. Van Nostrand Company, Inc., New York, 3rd edn., 1941.
21. D. Lide (ed.), *CRC Handbook of Chemistry and Physics, 86th Edition*, CRC Press, Boca Raton, FL, 2005.
22. J. Matic, X. Sun, B. A. Garetz and A. S. Myerson, *Cryst. Growth. Des.*, 2005, **5**, 1565-1567.
23. C. A. Schneider, W. S. Rasband and K. W. Eliceiri, *Nat. Methods*, 2012, **9**, 671-675.
24. D. J. Sheskin, *Handbook of Parametric and Nonparametric Statistical Procedures*, CRC Press, Boca Raton, FL, 5th edn., 2011.
25. X. Y. Sun, B. A. Garetz, M. F. Moreira and P. Palffy-Muhoray, *Phys. Rev. E*, 2009, **79**, 021710.

26. B. Clair, A. Ikni, W. Li, P. Scouflaire, V. Quemener and A. Spasojević-de Biré, *J. Appl. Crystallogr.*, 2014, **47**, 1252-1260.
27. N. Javid, T. Kendall, I. S. Burns and J. Sefcik, *Cryst. Growth. Des.*, 2016, **16**, 4196-4202.
28. B. C. Knott, J. L. LaRue, A. M. Wodtke, M. F. Doherty and B. Peters, *J. Chem. Phys.*, 2011, **134**, 171102.