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Observation of the Doubly Charmed Baryon Ξ_{cc}^{++}

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A highly significant structure is observed in the $\Lambda_c^+ K^- \pi^+ \pi^+$ mass spectrum, where the Λ_c^+ baryon is reconstructed in the decay mode $pK^-\pi^+$. The structure is consistent with originating from a weakly decaying particle, identified as the doubly charmed baryon Ξ_{cc}^{++}. The difference between the masses of the Ξ_{cc}^{++} and Λ_c^+ states is measured to be $1334.94 \pm 0.72(\mathrm{stat.}) \pm 0.27(\mathrm{syst.})\ \mathrm{MeV}/c^2$, and the Ξ_{cc}^{++} mass is then determined to be $3621.40 \pm 0.72(\mathrm{stat.}) \pm 0.27(\mathrm{syst.}) \pm 0.14(\Lambda_c^+)\ \mathrm{MeV}/c^2$, where the last uncertainty is due to the limited knowledge of the Λ_c^+ mass. The state is observed in a sample of proton-proton collision data collected by the LHCb experiment at a center-of-mass energy of 13 TeV, corresponding to an integrated luminosity of $1.7\ \mathrm{fb}^{-1}$, and confirmed in an additional sample of data collected at 8 TeV.

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The quark model [1–3] predicts the existence of multiplets of baryon and meson states. Those states composed of the lightest four quarks (u, d, s, c) form SU(4) multiplets [4]. Numerous states with charm quantum number $C = 0$ or $C = 1$ have been discovered, including all of the expected $q\bar{q}$ and qqq ground states [5]. Three weakly decaying qqq states with $C = 2$ are expected: one isospin doublet ($\Xi_{cc}^{++} = ccu$ and $\Xi_{cc}^{+} = ccd$) and one isospin singlet ($\Omega_{cc}^{+} = ccs$), each with spin parity $J^P = 1/2^+$. The properties of these baryons have been calculated with a variety of theoretical models. In most cases, the masses of the Ξ_{cc} states are predicted to lie in the range 3500 to 3700 MeV/c^2 [6–33]. The masses of the Ξ_{cc}^{++} and Ξ_{cc}^{+} states are expected to differ by only a few MeV/c^2, due to approximate isospin symmetry [34–36]. Most predictions for the lifetime of the Ξ_{cc}^{+} baryon are in the range 50 to 250 fs, and the lifetime of the Ξ_{cc}^{++} baryon is expected to be three to four times longer at 200 to 700 fs [10,11,19,24, 37–40]. While both are expected to be produced at hadron colliders [41–43], the longer lifetime of the Ξ_{cc}^{++} baryon should make it significantly easier to observe than the Ξ_{cc}^{+} baryon in such experiments, due to the use of real-time (online) event-selection requirements designed to reject backgrounds originating from the primary interaction point.

Experimentally, there is a long-standing puzzle in the Ξ_{cc} system. Observations of the Ξ_{cc}^{+} baryon at a mass of $3519 \pm 2\ \mathrm{MeV}/c^2$ with signal yields of 15.9 events over 6.1 ± 0.5 background in the final state $\Lambda_c^+ K^- \pi^+$ (6.3σ significance), and 5.62 events over 1.38 ± 0.13 background in the final state pD^+K^- (4.8σ significance) were reported by the SELEX Collaboration [44,45]. Their results included a number of unexpected features, notably a short lifetime and a large production rate relative to that of the singly charmed Λ_c^+ baryon. The lifetime was stated to be shorter than 33 fs at the 90% confidence level, and SELEX concluded that 20% of all Λ_c^+ baryons observed by the experiment originated from Ξ_{cc}^{+} decays, implying a relative Ξ_{cc} production rate several orders of magnitude larger than theoretical expectations [11]. Searches from the FOCUS [46], *BABAR* [47], and Belle [48] experiments did not find evidence for a state with the properties reported by SELEX, and neither did a search at LHCb with data collected in 2011 corresponding to an integrated luminosity of $0.65\ \mathrm{fb}^{-1}$ [49]. However, because the production environments at these experiments differ from that of SELEX, which studied collisions of a hyperon beam on fixed nuclear targets, these null results do not exclude the original observations.

This Letter presents the observation of the Ξ_{cc}^{++} baryon [50] via the decay mode $\Lambda_c^+ K^- \pi^+ \pi^+$ (Fig. 1), which is expected to have a branching fraction of up to 10% [51]. The Λ_c^+ baryon is reconstructed in the final state $pK^-\pi^+$.

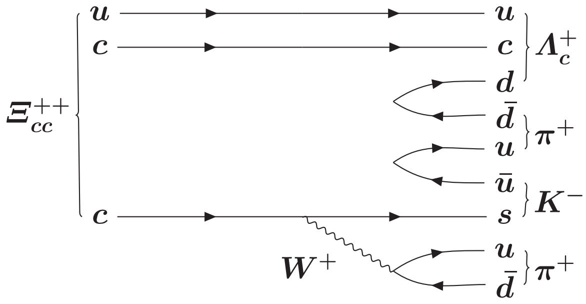

FIG. 1. Example Feynman diagram contributing to the decay $\Xi_{cc}^{++} \to \Lambda_c^+ K^- \pi^+ \pi^+$.

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The data consist of pp collisions collected by the LHCb experiment at the Large Hadron Collider at CERN with a center-of-mass energy of 13 TeV taken in 2016, corresponding to an integrated luminosity of 1.7 fb^{-1}.

The LHCb detector is a single-arm forward spectrometer covering the pseudorapidity range $2 < \eta < 5$, designed for the study of particles containing b or c quarks, and is described in detail in Refs. [52,53]. The detector elements most relevant to this analysis are a silicon-strip vertex detector surrounding the pp interaction region, a tracking system that provides a measurement of the momentum of charged particles, and two ring-imaging Cherenkov detectors [54] that are able to discriminate between different species of charged hadrons. The on-line event selection is performed by a trigger that consists of a hardware stage, which is based on information from the calorimeter and muon systems, followed by a software stage, which fully reconstructs the event [55]. The on-line reconstruction incorporates near-real-time alignment and calibration of the detector [56], which in turn allows the reconstruction of the Ξ_{cc}^{++} decay to be performed entirely in the trigger software.

The reconstruction of $\Xi_{cc}^{++} \to \Lambda_c^+ K^- \pi^+ \pi^+$ decays proceeds as follows. Candidate $\Lambda_c^+ \to p K^- \pi^+$ decays are reconstructed from three charged particles that form a good-quality vertex and that are inconsistent with originating from any pp collision primary vertex (PV). The PV of any single particle is defined to be the PV with respect to which the particle has the smallest impact parameter χ^2 (χ^2_{IP}), which is the difference in χ^2 of the PV fit with and without the particle in question. The Λ_c^+ vertex is required to be displaced from its PV by a distance corresponding to a proper decay time greater than 150 fs. The Λ_c^+ candidate is then combined with three additional charged particles to form a $\Xi_{cc}^{++} \to \Lambda_c^+ K^- \pi^+ \pi^+$ candidate. These additional particles must form a good-quality vertex with the Λ_c^+ candidate, and the Λ_c^+ decay vertex must be downstream of the Ξ_{cc}^{++} vertex. Each of the six final-state particles is required to pass track-quality requirements, to have hadron-identification information consistent with the appropriate hypothesis (p, K, or π), and to have transverse momentum $p_T > 500\ \mathrm{MeV}/c$. To avoid duplicate tracks, the angle between each pair of final-state particles with the same charge is required to be larger than 0.5 mrad. The Ξ_{cc}^{++} candidate must have $p_T > 4\ \mathrm{GeV}/c$ and must be consistent with originating from its PV. The selection above includes criteria applied in the trigger software, plus additional requirements chosen based on simulated signal events and a control sample of data. Simulated signal events are produced with the standard LHCb simulation software [57–63] interfaced to a dedicated generator, GENXICC [64–66], for Ξ_{cc}^{++} baryon production. In the simulation, the Ξ_{cc}^{++} mass and lifetime are assumed to be $3600\ \mathrm{MeV}/c^2$ and 333 fs. The background control sample consists of wrong-sign (WS) $\Lambda_c^+ K^- \pi^+ \pi^-$ combinations.

The background level is further reduced with a multivariate selector based on the multilayer perceptron algorithm [67]. The selector is trained with simulated signal events and with the WS control sample of data to represent the background. For both signal and background training samples, candidates are required to pass the selection described above and to fall within a signal search region defined as $2270 < m_{\mathrm{cand}}(\Lambda_c^+) < 2306\ \mathrm{MeV}/c^2$ and $3300 < m_{\mathrm{cand}}(\Xi_{cc}^{++}) < 3800\ \mathrm{MeV}/c^2$, where $m_{\mathrm{cand}}(\Lambda_c^+)$ is the reconstructed mass of the Λ_c^+ candidate, $m_{\mathrm{cand}}(\Xi_{cc}^{++}) \equiv m(\Lambda_c^+ K^- \pi^+ \pi^\pm) - m_{\mathrm{cand}}(\Lambda_c^+) + m_{\mathrm{PDG}}(\Lambda_c^+)$, $m(\Lambda_c^+ K^- \pi^+ \pi^\pm)$ is the reconstructed mass of the $\Lambda_c^+ K^- \pi^+ \pi^\pm$ combination, and $m_{\mathrm{PDG}}(\Lambda_c^+) = 2286.46 \pm 0.14\ \mathrm{MeV}/c^2$ is the known value of the Λ_c^+ mass [5]. The $m_{\mathrm{cand}}(\Lambda_c^+)$ window corresponds to approximately ± 3 times the Λ_c^+ mass resolution. The use of $m_{\mathrm{cand}}(\Xi_{cc}^{++})$ rather than $m(\Lambda_c^+ K^- \pi^+ \pi^\pm)$ cancels fluctuations in the reconstructed Λ_c^+ mass to first order, and thereby improves the Ξ_{cc}^{++} mass resolution by approximately 40%.

Based on studies with simulated events and control samples of data, ten input variables that together provide good discrimination between signal and background candidates are used in the multivariate selector. They are as follows: the χ^2 per degree of freedom of each of the Λ_c^+ vertex fit, the Ξ_{cc}^{++} vertex fit, and a kinematic refit [68] of the Ξ_{cc}^{++} decay chain requiring it to originate from its PV; the smallest p_T of the three decay products of the Λ_c^+; the smallest p_T of the four decay products of the Ξ_{cc}^{++}; the scalar sum of the p_T of the four decay products of the Ξ_{cc}^{++}; the angle between the Ξ_{cc}^{++} momentum vector and the direction from the PV to the Ξ_{cc}^{++} decay vertex; the flight distance χ^2 between the PV and the Ξ_{cc}^{++} decay vertex; the χ^2_{IP} of the Ξ_{cc}^{++} with respect to its PV; and the smallest χ^2_{IP} of the decay products of the Ξ_{cc}^{++} with respect to its PV. Here, the flight distance χ^2 is defined as the χ^2 of the hypothesis that the Ξ_{cc}^{++} decay vertex coincides with its PV. Candidates are retained for analysis only if their multivariate selector output values exceed a threshold chosen by maximizing the expected value of the figure of merit $\varepsilon/(\frac{5}{2} + \sqrt{B})$ [69], where ε is the estimated signal efficiency and B is the estimated number of background candidates underneath the signal peak. The quantity B is computed with the WS control sample and, purely for the purposes of this optimization, it is calculated in a window centered at a mass of $3600\ \mathrm{MeV}/c^2$ and of half-width $12.5\ \mathrm{MeV}/c^2$ (corresponding to approximately twice the expected resolution). Its evaluation takes into account the difference in background rates between the $\Lambda_c^+ K^- \pi^+ \pi^+$ signal mode and the WS sample, scaling the WS background by the ratio seen in data in the sideband regions $3200 < m_{\mathrm{cand}}(\Xi_{cc}^{++}) < 3300\ \mathrm{MeV}/c^2$ and $3800 < m_{\mathrm{cand}}(\Xi_{cc}^{++}) < 3900\ \mathrm{MeV}/c^2$. The performance of the multivariate selector is also tested for simulated signal events under other lifetime hypotheses; while the signal efficiency increases with the lifetime, it is

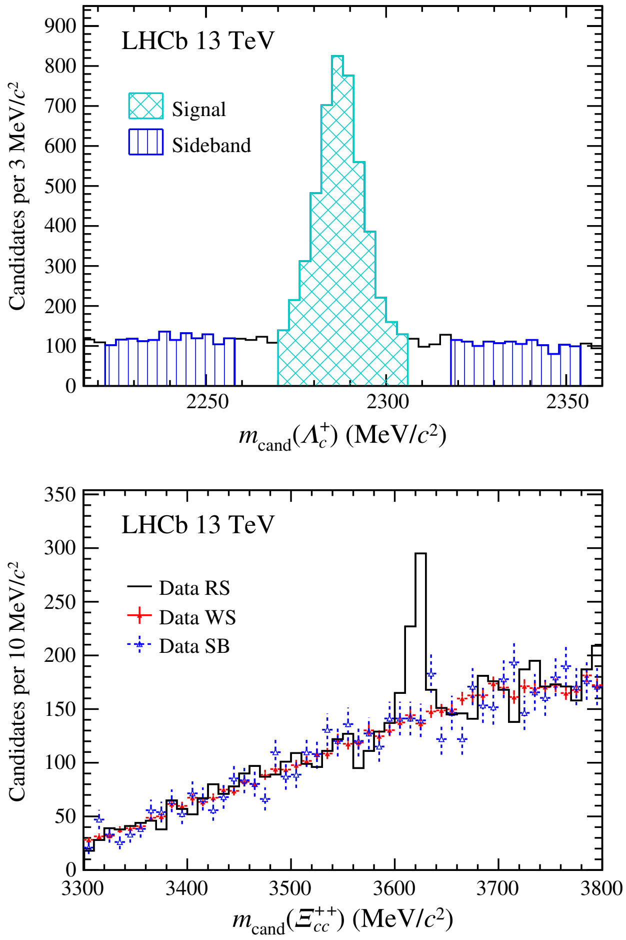

FIG. 2. Mass spectra of (upper) Λ_c^+ and (lower) Ξ_{cc}^{++} candidates. The full selection is applied, except for the Λ_c^+ mass requirement in the case of the upper plot. For the Λ_c^+ mass distribution the (cross-hatched) signal and (vertical line) sideband regions are indicated; to avoid duplication, the histogram is filled only once in events that contain more than one Ξ_{cc}^{++} candidate. In the lower plot the right-sign (RS) signal sample $\Xi_{cc}^{++} \to \Lambda_c^+ K^- \pi^+ \pi^+$ is shown, along with the control samples: Λ_c^+ sideband (SB) $\Lambda_c^+ K^- \pi^+ \pi^+$ candidates and wrong-sign (WS) $\Lambda_c^+ K^- \pi^+ \pi^-$ candidates, normalized to have the same area as the RS sample in the $m_{\text{cand}}(\Xi_{cc}^{++})$ sidebands.

found that the training obtained for 333 fs is close to optimal (i.e., gives comparable performance to a training optimized for the new lifetime hypothesis) even for much shorter or longer lifetimes.

After the multivariate selection is applied, events may still contain more than one Ξ_{cc}^{++} candidate in the signal search region. Based on studies of simulation and the control data sample, no peaking background arises due to multiple candidates except for the special case in which the candidates are formed from the same six decay products but two of the decay products are interchanged (e.g., the K^- particle from the Ξ_{cc}^{++} decay and the K^- particle from the Λ_c^+ decay). In such instances, one of the candidates is chosen at random to be retained and all others are discarded. In the remaining events, the fraction that has more than one Ξ_{cc}^{++} candidate in the range 3300–3800 MeV/c^2 is approximately 8%.

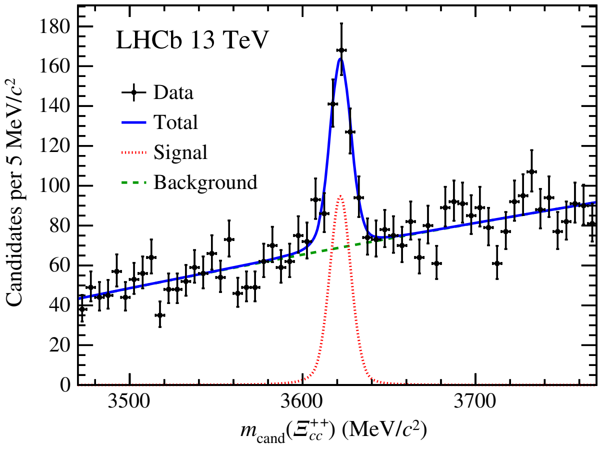

FIG. 3. Invariant mass distribution of $\Lambda_c^+ K^- \pi^+ \pi^+$ candidates with fit projections overlaid.

The selection described above is then applied to data in the search region. Figure 2 shows the Λ_c^+ mass distribution, and the Ξ_{cc}^{++} mass spectra for candidates in the mass range $2270 < m_{\text{cand}}(\Lambda_c^+) < 2306$ MeV/c^2. A structure is visible in the signal mode at a mass of approximately 3620 MeV/c^2. No significant structure is visible in the WS control sample, or for events in the Λ_c^+ mass sidebands. To measure the properties of the structure, an unbinned extended maximum likelihood fit is performed to the invariant mass distribution in the restricted $\Lambda_c^+ K^- \pi^+ \pi^+$ mass window of 3620 ± 150 MeV/c^2 (Fig. 3). The peaking structure is empirically described by a Gaussian function plus a modified Gaussian function with power-law tails on both sides [70]. All peak parameters are fixed to values obtained from simulation apart from the mass, yield, and an overall resolution parameter. The background is described by a second-order polynomial with parameters free to float in the fit. The signal yield is measured to be 313 ± 33, corresponding to a local statistical significance in excess of 12σ when evaluated with a likelihood ratio test. The fitted resolution parameter is 6.6 ± 0.8 MeV/c^2, consistent with simulation. The same structure is also observed in the $\Lambda_c^+ K^- \pi^+ \pi^+$ spectrum in a pp data sample collected by LHCb at $\sqrt{s} = 8$ TeV (see the Supplemental Material [71] for results from the 8 TeV cross-check sample). The local statistical significance of the peak in the 8 TeV sample is above 7σ, and its mass is consistent with that in the 13 TeV data sample.

Additional cross-checks are performed confirming the robustness of the observation. The significance of the structure in the $\Lambda_c^+ K^- \pi^+ \pi^+$ final state remains above 12σ when fixing the resolution parameter in the invariant mass fit to the value obtained from simulation, changing the threshold value for the multivariate selector, removing events containing multiple candidates in the fitted mass

TABLE I. Systematic uncertainties on the Ξ_{cc}^{++} mass measurement.

Source	Value (MeV/c^2)
Momentum-scale calibration	0.22
Selection bias correction	0.14
Unknown Ξ_{cc}^{++} lifetime	0.06
Mass fit model	0.07
Sum of above in quadrature	0.27
Λ_c^+ mass uncertainty	0.14

range, or using an alternative selection without a multivariate classifier. The significance also remains above 12σ in a subsample of candidates for which the reconstructed decay time exceeds five times its uncertainty. This is consistent with a weakly decaying state and inconsistent with the strong decay of a resonance. No fake peaking structures are observed in the control samples when requiring various intermediate resonances to be present (ρ^0, K^{*0}, Σ_c^0, Σ_c^{++}, Λ_c^{*+}) nor are they observed when combining Ξ_{cc}^{++} and Λ_c^+ decay products. The contributions of misidentified $D_s^+ \to K^+K^-\pi^+$ and $D^+ \to K^-\pi^+\pi^+$ decays are found to be negligible.

The sources of systematic uncertainty affecting the measurement of the Ξ_{cc}^{++} mass (Table I) include the momentum-scale calibration, the event selection, the unknown Ξ_{cc}^{++} lifetime, the invariant mass fit model, and the uncertainty on the Λ_c^+ mass. The momentum scale is calibrated with samples of $J/\psi \to \mu^+\mu^-$ and $B^+ \to J/\psi K^+$ decays [72,73]. After calibration, an uncertainty of ±0.03% is assigned, which corresponds to a systematic uncertainty of 0.22 MeV/c^2 on the reconstructed Ξ_{cc}^{++} mass. The selection procedure is more efficient for vertices that are well separated from the PV, and therefore preferentially retains longer-lived Ξ_{cc}^{++} candidates. Because of a correlation between the reconstructed decay time and the reconstructed mass, this induces a positive bias on the mass for both Ξ_{cc}^{++} and Λ_c^+ candidates. The effect is studied with simulation and the bias on the Ξ_{cc}^{++} mass is determined to be $+0.45 \pm 0.14$ MeV/c^2 (assuming a lifetime of 333 fs), where the uncertainty is due to the limited size of the simulation sample. A corresponding correction is applied to the fitted value in data. To validate this procedure, the Λ_c^+ mass in an inclusive sample is measured and corrected in the same way; after the correction, the Λ_c^+ mass is found to agree with the known value [5]. The bias on the Ξ_{cc}^{++} mass depends on the unknown Ξ_{cc}^{++} lifetime, introducing a further source of uncertainty on the correction. This is estimated by repeating the procedure for other Ξ_{cc}^{++} lifetime hypotheses between 200 and 700 fs. The largest deviation in the correction, 0.06 MeV/c^2, is taken as an additional systematic uncertainty. Final-state photon radiation also causes a bias in the measured mass, which is determined to be -0.05 MeV/c^2 with simulation [61]. The uncertainty on this correction is approximately 0.01 MeV/c^2 and is neglected. The dependence of the measurement on the fit model is estimated by varying the shape parameters that are fixed according to simulation, by using alternative signal and background models, and by repeating the fits in different mass ranges. The largest deviation seen in the mass, 0.07 MeV/c^2, is assigned as a systematic uncertainty. Finally, since the Ξ_{cc}^{++} mass is measured relative to the Λ_c^+ mass, the uncertainty of 0.14 MeV/c^2 on the world-average value of the latter is included. After taking these systematic effects into account and combining their uncertainties (except that on the Λ_c^+ mass) in quadrature, the Ξ_{cc}^{++} mass is measured to be $3621.40 \pm 0.72(\text{stat.}) \pm 0.27(\text{syst.}) \pm 0.14(\Lambda_c^+)$ MeV/c^2. The mass difference between the Ξ_{cc}^{++} and Λ_c^+ states is $1334.94 \pm 0.72(\text{stat.}) \pm 0.27(\text{syst.})$ MeV/c^2.

In summary, a highly significant structure is observed in the final state $\Lambda_c^+ K^-\pi^+\pi^+$ in a pp data sample collected by LHCb at $\sqrt{s} = 13$ TeV, with a signal yield of 313 ± 33. The mass of the structure is measured to be $3621.40 \pm 0.72(\text{stat.}) \pm 0.27(\text{syst.}) \pm 0.14(\Lambda_c^+)$ MeV/c^2, where the last uncertainty is due to the limited knowledge of the Λ_c^+ mass, and its width is consistent with experimental resolution. The structure is confirmed with consistent mass in a data set collected by LHCb at $\sqrt{s} = 8$ TeV. The signal candidates have significant decay lengths, and the signal remains highly significant after a minimum lifetime requirement of approximately five times the expected decay-time resolution is imposed. This state is therefore incompatible with a strongly decaying particle but is consistent with the expectations for the weakly decaying Ξ_{cc}^{++} baryon. The mass of the observed Ξ_{cc}^{++} state is greater than that of the Ξ_{cc}^{+} peaks reported by the SELEX Collaboration [44,45] by 103 ± 2 MeV/c^2. This difference would imply an isospin splitting vastly larger than that seen in any other baryon system and is inconsistent with the expected size of a few MeV/c^2 [34–36]. Consequently, while the state reported here is consistent with most theoretical expectations for the Ξ_{cc}^{++} baryon, it is inconsistent with being an isospin partner to the Ξ_{cc}^{+} state reported previously by the SELEX Collaboration.

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